

Discovery and Mass Spectroscopy via Mixed Di-Boson Production at the LHC in a 4-Site Model with a Composite Higgs Boson

Luca Fedeli* and Stefano Moretti*

*School of Physics and Astronomy, University of Southampton,
Highfield, Southampton SO17 1BJ, UK*

Stefania De Curtis†

*Istituto Nazionale di Fisica Nucleare, Sezione di Firenze,
Via G. Sansone 1, 50019 Sesto Fiorentino, Italy*

Daniele Dominici‡

*Università degli Studi di Firenze, Dip. di Fisica e Astronomia,
and Istituto Nazionale di Fisica Nucleare, Sezione di Firenze,
Via G. Sansone 1, 50019 Sesto Fiorentino, Italy*

(Dated: October 30, 2012)

We investigate the scope of the Large Hadron Collider (LHC) in probing the parameter space of a 4-Site (4S) model supplemented by one composite Higgs state. We concentrate on the yield of ‘mixed’ di-boson production giving three leptons and missing (transverse) energy in the final state, i.e., events induced via the subprocess $q\bar{q}' \rightarrow l^+\nu_l l'^+l'^-$ + c.c., where q, q' are quarks and $l, l' = e, \mu$ in all possible combinations, signature which enables the production in the intermediate steps of all additional – with respect to the Standard Model (SM) – neutral and charged gauge bosons belonging to the spectrum of this scenario, all of which in resonant topologies. We not only find this channel to be accessible over the background but also show that, after a dedicated cut-based analysis, kinematical reconstruction of all such resonances is always possible. We finally compare the yield of the mixed di-boson mode to that of the ‘charged’ di-boson one and find that the former is of lesser scope than the latter at smaller gauge boson masses but it overcomes it at higher values, thereby revealing itself a viable search mode covering new regions of parameter space.

PACS numbers: 12.60.Cn, 11.25.Mj, 12.39.Fe

I. INTRODUCTION

A strongly broken Electro-Weak (EW) sector is expected to contain a variety of bound states including particles of spin 0, spin 1/2 as well as spin 1. Its phenomenology below the scale of the new interactions responsible for EW Symmetry Breaking (EWSB) and producing such new resonances can be studied, amongst others, in simple four-dimensional models, like the 4-Site (4S) one of [1].

In its original formulation, the 4S model describes in an effective way the interactions of extra spin 1 resonances (four in total) as gauge fields of a $SU(2) \otimes SU(2)$ extra gauge group. They can be thought of as the first Kaluza-Klein (KK) excitations emerging from a five-dimensional formulation, and, due to the Anti-de Sitter/Conformal Field Theory (AdS/CFT) correspondence, they are composite states of a strong dynamics also responsible for the breaking of the EW symmetry. In this note we consider a 4S model supplemented by the inclusion of a new scalar field, singlet under the above gauge group, in order to reproduce, in our effective description, the scalar particle recently detected by the ATLAS and CMS experiments at the Large Hadron Collider (LHC) [2, 3]. The model has been introduced in [4], to which we refer the reader for technical details. In such a scenario, the couplings of our composite Higgs-like scalar particle to the Standard Model (SM) and extra gauge bosons are free parameters for which we have derived bounds due to the EW Precision Tests (EWPTs) and the present measurements at the Tevatron and LHC, via Drell-Yan (DY) processes, as well as theoretical constraints enforced by unitarity requirements.

It is the purpose of this paper to investigate, in the context of such a 4S model with one composite Higgs state (compatible with Tevatron and LHC data), the phenomenology of ‘mixed’ di-boson production at the LHC, yielding three leptons and missing transverse energy, i.e., the process

$$pp(q\bar{q}') \rightarrow W^\pm Z \rightarrow l^\pm l'^+ l'^- E_T^{\text{miss}}, \quad (1)$$

wherein the symbols W^\pm and Z refer to any possible charged and neutral, respectively, spin 1 massive gauge bosons present in the model, q and q' are the quarks found inside the proton whilst l and l' refer to either an electron e or a muon μ , the latter two in all possible

^{*}Electronic address: l.fedeli@soton.ac.uk; s.moretti@soton.ac.uk

[†]Electronic address: decurtis@fi.infn.it

[‡]Electronic address: dominici@fi.infn.it

combinations of (identical and different) such flavours.

In the attempt to extract the signal in (1), we will focus our attention onto all energy and luminosity stages covered already or still foreseen for the CERN machine. Ultimately, we will want to contrast the detection potential at the LHC of mixed di-boson production with that of both DY,

$$q\bar{q}' \rightarrow W^\pm \rightarrow l^\pm E_T^{\text{miss}}, \quad q\bar{q} \rightarrow Z \rightarrow l^+l^-, \quad (2)$$

and ‘charged’ di-boson events,

$$pp(q\bar{q}) \rightarrow W^+W^- \rightarrow e^\pm\mu^\mp E_T^{\text{miss}}, \quad (3)$$

building on previous studies of ours. Further, we will show that the presence of only one source of missing transverse energy (the neutrino ν_ℓ) in the case of mixed di-boson production enables one to extract simultaneously, upon reconstructing its longitudinal momentum, all masses of the new gauge bosons present in the model.

The plan of the paper is as follows. The next section will describe the model at hand. The following one, after describing the calculation, will be devoted to study the production and decay dynamics of process (1), eventually extracting from it exclusion and discovery limits over the accessible parameter space and contrasting these to the yield of processes (2) and (3). Then, in the same section, we will also illustrate how to extract the mass spectrum of the 4S model. A final section will be devoted to summarise our work and conclude.

II. THE 4S MODEL WITH A SINGLET COMPOSITE SCALAR HIGGS STATE

The 4S model is a moose model based on a $SU(2)_L \otimes SU(2)_1 \otimes SU(2)_2 \otimes U(1)_Y$ gauge symmetry and contains three non-linear σ -model fields interacting with the gauge fields, which trigger spontaneous EWSB. Its construction is presented in [1, 4] while some of its phenomenological consequences are analysed in [5, 6, 7, 8, 9, 10]. In the unitary gauge, the spectrum of the 4S model incorporates (other than all the SM particles, including the Higgs boson, which is however not responsible for EWSB) two new triplets of gauge bosons, which acquire mass through the same non-linear symmetry breaking mechanism giving mass to the SM gauge bosons. We denote with $W_{i\mu}^\pm$ and $Z_{i\mu}$ ($i = 1, 2$) the four charged and two neutral heavy resonances, respectively, appearing as a consequence of the gauge group extension,

and with W_μ^\pm , Z_μ and A_μ the SM gauge bosons. Owing to its gauge structure, the 4S model a priori contains seven free parameters: the $SU(2)_L \otimes U(1)_Y$ gauge couplings, \tilde{g} and \tilde{g}' , the extra $SU(2)_{1,2}$ gauge couplings that we assume to be equal, $g_2 = g_1$, due to a LR symmetry [1], the bare masses of lighter (W_1^\pm, Z_1) and heavier (W_2^\pm, Z_2) gauge boson triplets, $M_{1,2}$, and their bare direct couplings to SM fermions, $b_{1,2}$, as described in [1]. However, their number can be reduced to four, by fixing the gauge couplings $\tilde{g}, \tilde{g}', g_1$ in terms of the three SM input parameters α, G_F, M_Z , which denote the Electro-Magnetic (EM) and Fermi constants plus the Z boson mass, respectively. As a result, the parameter space is completely defined by four independent free parameters, which one can choose to be: M_1, z, b_1 and b_2 , where $z = M_1/M_2$ is the ratio between the bare masses. In terms of these four parameters, physical masses and couplings of the extra gauge bosons to ordinary matter can be obtained via a complete numerical algorithm [8]. In the following, we choose to describe the full parameter space via the physical observables: other than z (which, as shown in [8], is a good approximation of the ratio between physical masses M_{W_1}/M_{W_2} or M_{Z_1}/M_{Z_2}), we take M_1, a_{W_1} and a_{W_2} which denote the mass of the lighter extra gauge boson and the couplings of the lighter and heavier extra gauge bosons to ordinary matter, respectively.

We refer the reader to Refs. [1, 4] for all details about theoretical and experimental bounds, mass spectrum and new resonance properties of the model being studied here.

III. RESULTS

A. Calculation

The codes exploited for our study of the LHC signatures are based on helicity amplitudes, defined through the HELAS subroutines [11] and assembled by means of MadGraph [12]. VEGAS [13] was used for the multi-dimensional numerical integrations. The Matrix Elements (ME) accounts for all off-shellness effects of the particles involved and was constructed starting from the topologies in Fig. 1, wherein (as intimated) the labels Z and W refer to any possible combination of gauge bosons in our model. The Parton Distribution Functions (PDFs) used were CTEQ5L [14], with factorisation/renormalisation scale set to $Q = \mu = \sqrt{s}$. Initial state quarks have been taken as massless, just like, unless otherwise stated, the final state leptons and neutrino.

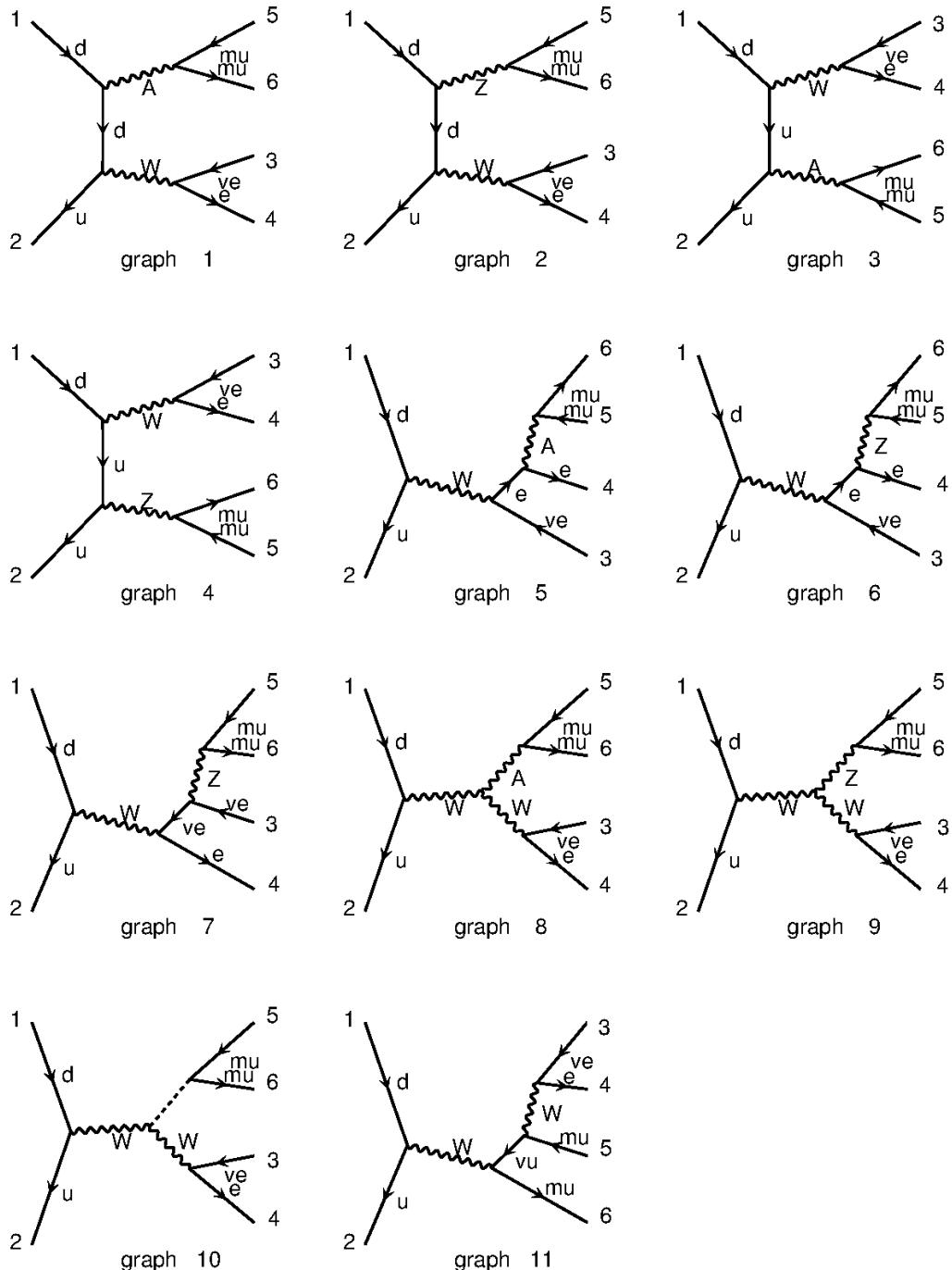


FIG. 1: Topologies of Feynman diagrams for the process in eq. (1), in the case $l \neq l'$ (in this case, e.g., $l = e$ and $l' = \mu$). For the case $l = l'$, one ought to swap helicities/momenta of two identical particles and antisymmetrise the diagrams (see text). Here the labels Z and W refer to all possible gauge boson states of the model, neutral and charged, respectively.

B. Event Selection

To calculate the Cross-Section (CS) at the LHC for the mixed di-boson channel we start considering general purpose cuts, which we call Cc cuts. The ensuing selection is designed in a such way to have three leptons in the final state, to produce a large signal CS, thereby usable for search purposes at low luminosities, and to give some hint about the typical mass scales of the new gauge bosons. This set of cuts is designed in the same spirit as those previously used to study processes (2) [7, 9] and (3) [4], essentially based on detector acceptance.

However, as intimated in the Introduction, the presence of only one neutrino in the final state of process (1) enables us to fully reconstruct the final state kinematics (see later on for technical details). This, as we will show, affords us then with the possibility of not only improving the discovery power of mixed di-boson production but also extracting the new gauge boson resonances. However, in order to do so, we will need to introduce a second, much tighter set of additional cuts, denoted by $C2$, which will make use of the partial knowledge acquired through the previous, looser selection of the mass scales of the new 4S gauge boson resonances, specifically, of the lightest ones. This set of cuts will then be used to both compare the exclusion/discovery potential of the LHC among the three channels (1)–(3) and to enable one to observe all six gauge boson resonances (SM and 4S ones). The drawback is that the $C2$ sample yields a smaller CS than the Cc one, so that higher luminosities are required.

Before illustrating the results of the analysis, we ought to define the observables that have been used to characterise this process. Here and in the following, we refer to the (visible) particles through their numerical labels $i, j = 4, 5, 6$, as listed in Fig. 1.

- $P_{i[j]}^T = \sqrt{(P_i^x + P_j^x)^2 + (P_i^y + P_j^y)^2}$ is the transverse momentum of a particle i [of a pair of particles ij]
- $p_M^T = \max(P_4^T, P_5^T, P_6^T)$ is the maximum amongst the transverse momenta of the particles
- $P_M^T = \max(P_{45}^T, P_{46}^T, P_{56}^T)$ is the maximum amongst the transverse momenta of all possible pairs of particles

- $M_{ij[k]} = \sqrt{(P_i + P_j[+P_k])^\mu (P_i + P_j[+P_k])_\mu}$ is the invariant mass of a pair[tern] of particles
- $M_{ij[k]}^T = \sqrt{(P_i^0 + P_j^0[+P_k^0])^2 - (P_i^x + P_j^x[+P_k^x])^2 - (P_i^y + P_j^y[+P_k^y])^2}$ is the transverse mass of a pair[tern] of particles
- $\eta_i = -\log \tan \frac{\theta_i}{2}$ is the pseudo-rapidity of a particle
- $\theta_{i[j]}$ is the angle between the beam axis and a particle[between two particles]
- $\cos \phi_{ij} = \frac{P_i^x P_j^x + P_i^y P_j^y}{P_i^T P_j^T}$ is the cosine of the relative (azimuthal) angle between two particles in the plane transverse to the beam
- $E_T^{\text{miss}} = \sqrt{(P_4^x + P_5^x + P_6^x)^2 + (P_4^y + P_5^y + P_6^y)^2}$ is the missing transverse energy (due to the neutrino escaping detection)

Here, P_i ($i = 1, \dots, 6$) are the momenta of the external particles entering process (1). In the following we list the two sets of cuts used, as detailed above.

- Cc : $P_i^T > 20 \text{ GeV}$, $M_{ij} > 20 \text{ GeV}$, $E_T^{\text{miss}} > 50 \text{ GeV}$, $|\eta_i| < 2$, $\cos \theta_{ij} < 0.9$;
- $C2$: Cc plus $M_{456} > 0.9 M_{W_1}$, $P_{ij}^T > 150 \text{ GeV}$, $p_M^T > 150 \text{ GeV}$, $\cos \phi_{56} < -0.5$.

Before proceeding with the kinematical analysis though, a couple of subtle points should be noted. The signatures we are considering are four, differing only for the flavour of the final state leptons, as we include the charge conjugation directly in the Monte Carlo (MC) code. In particular, the fermion flavors may be the following: ($3 = e$, $4 = 5 = \mu$), ($3 = \mu$, $4 = 5 = e$), ($3 = 4 = 5 = e$) and ($3 = 4 = 5 = \mu$). We refer to the first two cases (which give identical rates) as the electron-muon ($e\mu$) combination and for the last two (also identical) as the three lepton ($3l$) one. For the $e\mu$ case, in order to obtain the total CS, one straightforwardly compute the diagrams in Fig. 1. In contrast, for the $3l$ case, one must take into account the fact that there are two identical leptons in the final state (the third charged lepton differ from the other two for its electric charge), so we have to consider an antisymmetric final state. This duplicates the number of diagrams to compute by exchanging the helicities ($h_{4,6}$) and momenta ($P_{4,6}$) of the two identical leptons, so that, according to Pauli's statistic, one obtains:

$$\mathcal{M}_{3l}(\dots, P_4, h_4, \dots, P_6, h_6) = \frac{1}{\sqrt{2}} (\mathcal{M}_{e\mu}(\dots, P_4, h_4, \dots, P_6, h_6) - \mathcal{M}_{e\mu}(\dots, P_6, h_6, \dots, P_4, h_4)). \quad (4)$$

Computing the total CS for both $e\mu$ and $3l$ we found a small variation ($\sigma_{3l}/\sigma_{e\mu} = 0.90 - 0.85$) for both the SM and the full 4S. Further, we have verified that, after either set of cuts is implemented, all of the differential distributions are essentially identical over most of the kinematic ranges considered.

Another important point in our analysis is the capability to identify the two leptons coming from the neutral particle (A, Z, Z_1, Z_2). In the $e\mu$ case this is trivial, since the couple of same flavor leptons comes for sure from a neutral gauge boson. In the $3l$ scenario this is in principle ambiguous, as we cannot distinguish between the two same charge leptons (there are two identical particles in the final state, so it is clear that they are indistinguishable), i.e., the one coming from the W, W_1, W_2 and the other coming from the A, Z, Z_1, Z_2 . However if we consider P_{45}^T and P_{56}^T (the couple 4,6 is useless, due to the fact that particles 4 and 6 have the same charge and cannot come from a neutral boson) we get that almost in the 100% of the cases $P_M^T = P_{56}^T$ (in fact, this is false only in 0.8% of the cases): Fig. 2 illustrates this. Therefore, we are in a position to distinguish amongst leptons also in the case of identical flavours and entitled to apply cuts to each of these individually. The above statement is

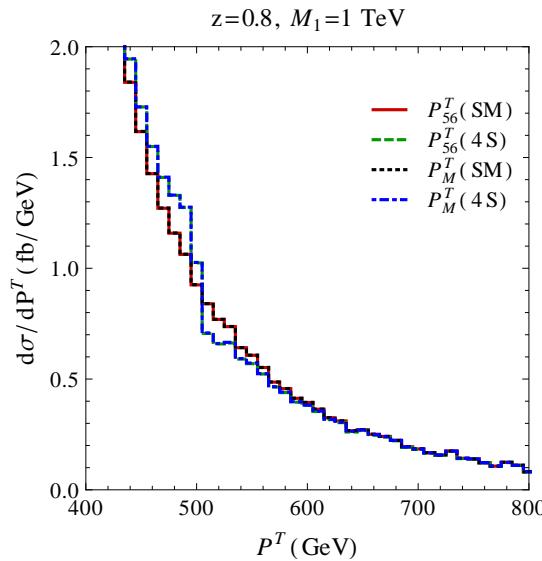


FIG. 2: Differential distributions for P_{56}^T and P_M^T for the SM and in a benchmark point of the 4S model (in particular, we choose $z = 0.8$, $M_1 = 1$ TeV and $a_{W_1} = 0.13$) after the Cc cuts. Notice that, on purpose, the lines are practically indistinguishable.

true after the Cc cuts, so that we can then exploit the above identification when applying

$\sqrt{s} = 7 \text{ TeV}$	$WW \text{ (fb)}$	$WZ \text{ (fb)}$	$\sqrt{s} = 8 \text{ TeV}$	$WW \text{ (fb)}$	$WZ \text{ (fb)}$	$\sqrt{s} = 14 \text{ TeV}$	$WW \text{ (fb)}$	$WZ \text{ (fb)}$
B	0.58	0.035	B	0.73	0.081	B	1.55	0.10
S	2.97	0.0003	S	3.99	0.0004	S	13.4	0.0013

TABLE I: Values of the CS for the 4S comparing processes (1), herein labeled WZ , and (3), herein labeled WW , for the case of a different flavor di-lepton signature, after the Cc cuts for the former and the So cuts (defined in Ref. [4]) for the latter.

the $C\mathcal{Q}$ cuts, i.e., those designed for mass spectrum extraction, for which the knowledge of the final state momenta is a prerequisite.

C. Mixed Di-Boson Production and Decay in the 4S Model

Before starting to consider the phenomenology of process (1) with three visible leptons, let us consider its contribution to the di-lepton sample exploited in Ref. [4], whereby the case of channel (3) was considered, limited to the different flavour case. Hence, for the process at hand here, we ought to consider its contribution when, e.g., the muon with the same charge of the electron is outside the detector acceptance region. In Tab. I we list the results using the Cc cuts for process (1) and the So cuts (defined in Ref. [4]) for channel (3), for past, present and future energy configurations of the LHC. Here, the benchmark point is defined by $z = 0.8$, $M_1 = 1 \text{ TeV}$ and $a_{W_1} = 0.13$ [16]. As we can see the WZ contribution is absolutely negligible compared to the WW one, to both the Background (B), which is the SM rate, and the Signal (S), defined as the difference between the full 4S rate and the SM one [17]. This is a general feature of the 4S model in fact, valid over its entire parameter space.

The situation is rather different for the three-lepton signatures, $e\mu$ and $3l$, that, from now on, we will treat cumulatively. In view of the fact that we will show that process (1) has some potential at the LHC in accessing the parameter space of the 4S model, it is worthed to review here the latter in some detail and explain how it can influence the CS for process (1). Of the aforementioned four independent parameters uniquely defining the 4S model, z , $M_1 (\sim M_{W_1} \sim M_{Z_1})$, a_{W_1} and a_{W_2} , we choose the first three as input parameters and fix the other in a such way to satisfy the EWPTs (namely $\epsilon_{1,2,3}$). This means that, for each set of

z , M_1 and a_{W_1} , there is a maximal ($a_{W_2}^M$) and a minimal ($a_{W_2}^m$) value allowed for a_{W_2} , e.g., for $z = 0.8$, $M_1 = 1$ TeV, $a_{W_1} = 0.13$ we have $-0.185 \leq a_{W_2} \leq -0.174$. To map the plane (M_{W_1}, a_{W_1}) we need to make a choice to fix a_{W_2} . Two possibilities are [18]:

$$(a) \quad a_{W_2}^a = \begin{cases} a_{W_2}^M & \text{if } |a_{W_2}^M| > |a_{W_2}^m| \\ a_{W_2}^m & \text{otherwise} \end{cases} \quad (5)$$

$$(b) \quad a_{W_2}^b = \begin{cases} a_{W_2}^M & \text{if } |a_{W_2}^M| < |a_{W_2}^m| \\ a_{W_2}^m & \text{otherwise} \end{cases} \quad (6)$$

In Tab. II we list some results for the case $z = 0.8$ and $M_1 = 1$ TeV. A choice made between cases (a) and (b) above influences also the other fermion couplings, for example the $f\bar{f}Z_2$ vertices. In particular, we get for the above combination of z and M_1 that (also assuming $a_{W_1} = 0.13$)

$$\frac{V_{f\bar{f}Z_2}^b}{V_{f\bar{f}Z_2}^a} \sim \frac{a_{W_2}^b}{a_{W_2}^a} = 0.94, \quad (7)$$

$$\frac{\Gamma_{W_2}^b}{\Gamma_{W_2}^a} \sim \frac{\Gamma_{Z_2}^b}{\Gamma_{Z_2}^a} \sim 0.97. \quad (8)$$

One could presume that the differences between choice (a) and (b) induced onto the CS

a_{W_1}	$a_{W_2}^a$	$a_{W_2}^b$	$a_{W_2}^b/a_{W_2}^a$
0.025	-0.022	-0.006	0.27
0.00	0.035	0.0018	0.51
-0.02	0.067	0.051	0.76

TABLE II: Values for a_{W_2} for three choices of a_{W_1} considering both cases (a) and (b) defined in eqs. (5) and (6), respectively, for $z = 0.8$ and $M_1 = 1$ TeV.

are negligible, especially when differences in masses and widths are rather small, as seen in eqs. (7)–(8). While this may be true in some cases, it is not so generally. Regarding the case of processes (2) and (3), this assumption has been verified to be correct in Refs. [4, 7, 9] (for the final choice of cuts used therein and imported here as well). For process (1) studied here, this is true for the $C2$ cuts (within typical uncertainties of CS measurements) but not the Cc ones, thereby further motivating the choice of the former as default set for our 4S parameter scans. In fact, assuming again $z = 0.8$ and $M_1 = 1$ TeV while taking, e.g., $a_{W_1} = 0.13$, we

get (using Cc cuts) $CS_a = 64$ fb and $CS_b = 88$ fb and (using $C2$ cuts) $CS_a = 1.8$ fb and $CS_b = 1.6$ fb. So it is clear that, on the one hand, for a given mass spectrum there is a large range available for the CS and, on the other hand, even if the mass is experimentally measured, there remains an uncertainty on the couplings. In essence, the numbers above already hint at the beneficial effects of the $C2$ cuts, with respect to the Cc ones, as the differences in CS between the choices in eqs. (5) and (6) are much less in the former case than in the latter. We will further quantify this effect later on in Subsect. III E. We now proceed though to compare channel (1) to processes (2) and (3) as discovery modes of the 4S model.

D. Exclusion and Discovery

We start this subsection by comparing the twin processes of mixed and charged di-boson production, i.e., (1) and (3), respectively, e.g., at 14 TeV. At inclusive level, as we can see from Tab. III, the main difference between the former (denoted by WZ) and the latter (denoted by WW) is a significantly different CS, notably due to a residual SM contribution entering process (1) with resonant dynamics, which is instead absent in process (3). However, the significance is better for the WW case than for the WZ one, primarily due to the aforementioned higher background induced by the SM in the latter case with respect to the former. Notice that, as previously, the background is the SM yield whilst the signal is the difference between the full 4S result and the SM itself. Similar results are obtained also at 7 and 8 TeV. Finally, such a pattern generally persists over the entire parameter space of the 4S model. If one instead adopts the $C2$ cuts, the WW and WZ rates become more compatible and the significances of the latter improve, to the expense of lower signal rates. Hence, it makes sense to adopt the tighter set of cuts as the default one in the parameter scan, further considering that they are designed to exalt the resonant contributions of the new gauge bosons states (unlike the looser set, which reveals a stronger sensitivity to interference effects), hence in tune with standard experimental approaches.

In the remainder of this section, we summarise the exclusion and discovery potential of the 14 TeV LHC with respect to processes (1)–(3). For the last two channels, we borrow some of the results obtained in Ref. [4]. Herein, for illustration purposes, we consider again the value of $z = 0.8$. Assuming $C2$ cuts and $e\mu$ final states for processes (1)–(3), the CS

M_1 (TeV)	0.5	0.75	1	1.5	1.7	2
WW	9.6	11.8	13.4	11.8	9.9	4.9
WZ (Cc)	57.0/33.4	89.2/64.6	89.0/63.8	60.8/41.4	56.8/36.4	32.0/51.2
WZ ($C2$)	1.70/1.42	1.42/1.20	1.80/1.56	1.82/1.68	1.50/1.36	1.08/0.98

TABLE III: Signal CS (in fb) defined as the total 4S rate minus the SM one at the LHC with $\sqrt{s} = 14$ TeV considering both di-boson processes discussed in the text, for various values of M_1 and z fixed to 0.8, with the choice of maximal allowed value for a_{W_1} . For the WZ process the two values reported here correspond to the solution (a)/(b) in eq. (5)/(6), for both the Cc and $C2$ cuts. The value for the SM is 1.55(5006/1.40) fb for the $WW(WZ[Cc/C2])$ channel. For both signatures, the di-lepton one emerging from WW and the tri-lepton one stemming from WZ , we have taken different flavours only, $e\mu$, for consistency.

decreases very quickly for a_{W_1} going to 0, consequently, there are only very small excluded regions through WZ in the (M_{W_1}, a_{W_1}) plane. In particular, we get that the WZ exclusion limits are always less stringent than those from WW . For what concerns discovery, the results are quite similar. This is illustrated in Fig. 3, where we show the exclusion and discovery limits at the LHC with 14 TeV after 15 fb^{-1} for processes (1) and (3), plotted against the regions still allowed by EWPTs and direct searches via the processes in (2), the latter updated to the latest CMS and ATLAS analyses based on 7 TeV data after 5 fb^{-1} , mapped over the (M_{W_1}, a_{W_1}) plane compliant with unitarity requirements. The above conclusions change though if, instead of solely relying on the $e\mu$ final state in the case of the WZ process, we also include the $3l$ case [19]. The same figure in fact also points out that process (1) equals or overtakes channel (3) (and in turn, also the modes in (2), see Ref. [4]), for exclusion and discovery purposes, respectively, at large gauge boson masses, yet still compatible with unitarity limits. Beyond these, for larger boson masses, because of kinematical reasons, the exclusion and discovery regions will eventually close.

Finally, notice that we have used here the choice (a) for a_{W_2} , though results are qualitatively the same for the case (b). In fact, even assuming 7 or 8 TeV for the LHC (and standard accrued luminosities of 5 and 15 fb^{-1} , respectively), the overall pattern remains unchanged though quantitatively the scope of the WZ mode is reduced in comparison to

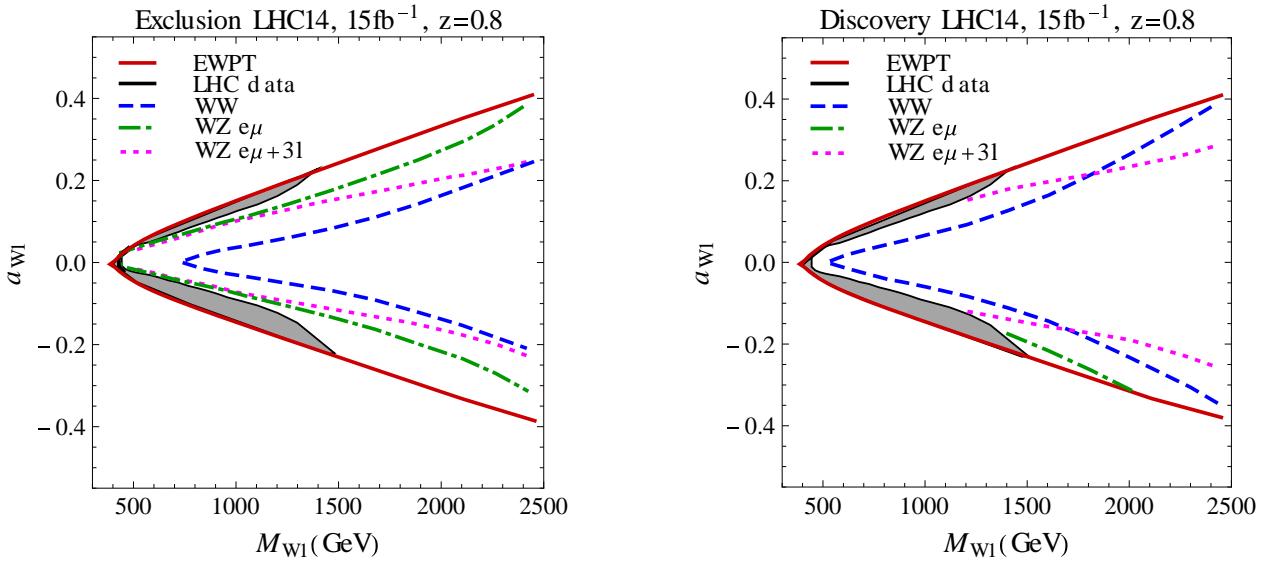


FIG. 3: Left(Right): Contrasting the WZ and WW yields in terms of 2σ exclusion(5σ discovery) contours over the plane (M_{W_1}, a_{W_1}) still allowed by EWPTs and direct searches via the DY modes for the indicated luminosities. We are using here the choice (*a*) for a_{W_2} . Here, we consider both the $e\mu$ and $3l$ final state in case of the WZ process. Also notice that the same detector efficiency for leptons and missing transverse energy as extracted from data (see Ref. [4] for further details) was used in all cases.

that of the WW and DY ones, owing to the fact that high gauge boson masses are more unattainable at reduced energies and luminosities.

E. Mass Spectroscopy

In this part of the paper, we intend to show that mixed di-boson production at the LHC (again, we take $\sqrt{s} = 14$ TeV for illustration purposes) can act as an effective means to extract the entire mass spectrum of the gauge bosons of the 4S model as well as to strongly constrain the size of their couplings. In order to accomplish this, it is crucial the fact that process (1) affords one with the possibility to reconstruct the missing transverse momentum of the neutrino. In particular, in order to do so, we use the algorithm of Ref. [15], which reproduces a (partonic) invariant mass ($\sqrt{\hat{s}} = E_{cm}$) distribution that agrees with the true distribution very well: see Fig. 4 [20]. Incidentally, the left-hand plot here highlights the

W_1 contribution, emerging from the interference with the SM, while the right-hand plot singles out the W_2 contribution, stemming from the 4S resonance. In principle, one should use all the three values of the intervening charged gauge boson masses (i.e., W , W_1 and W_2) to reconstruct the ν momentum. However, here we use only the SM mass, M_W , under the assumption that the corresponding resonance always gives a very sizable contribution amongst the three. In fact, this represents an effective choice, for both sets of cuts adopted throughout, i.e., Cc and $C2$ [21].

Again, the benchmark chosen here is representative of a situation occurring generally over the 4S parameter space.

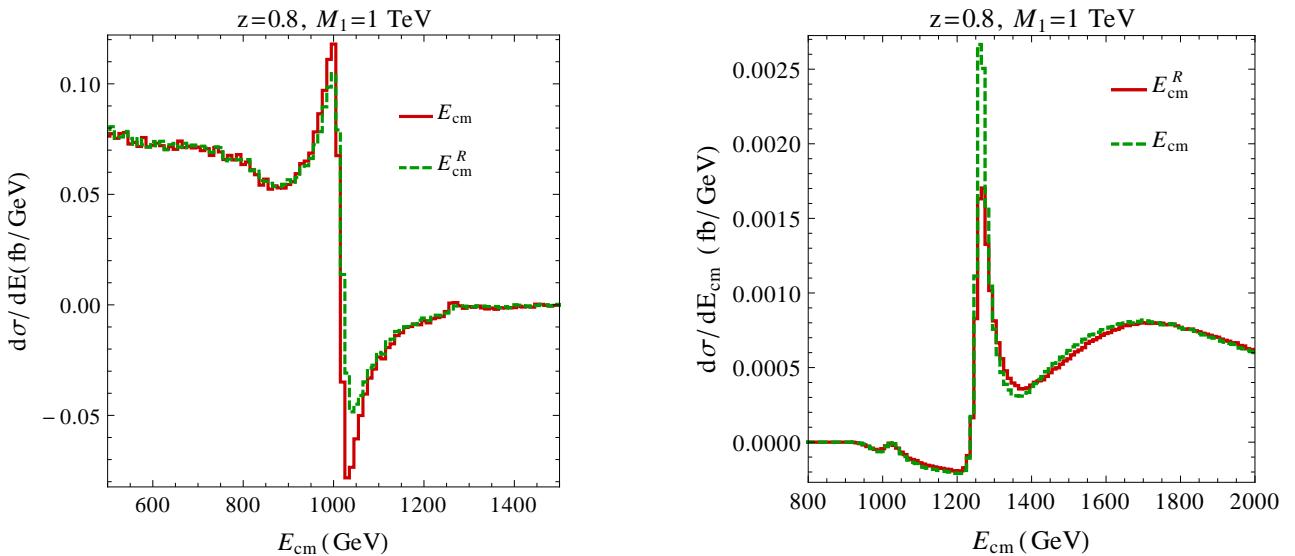


FIG. 4: Differential distributions of the true (E_{cm}) and reconstructed (E_{cm}^R) Center-of-Mass (CM) energy at partonic level for the signal, defined as the difference between the full 4S result and the SM one, for a benchmark point in parameter space. In particular, we choose $z = 0.8$, $M_1 = 1$ TeV and $a_{W_1} = 0.13$. Left(Right): using the $Cc(C2)$ cuts.

In Fig. 5 we show the differential distributions in M_{56} and E_{cm}^R (the ‘reconstructed’ E_{cm}), for a benchmark configuration of the 4S model ($z = 0.8$, $M_1 = 1$ TeV, $a_{W_1} = 0.13$). These two spectra are the most sensitive ones to the $Z_{1,2}$ and $W_{1,2}$ masses, respectively. The combination of the two set of cuts is necessary in order to extract the gauge boson mass spectrum of the 4S model. On the one hand, the Cc cuts can extract both neutral gauge boson resonances (top-left plot). For these in fact the $C2$ set merely acts in the direction of

further suppressing the SM background (notice the disappearance of the low invariant mass tail in the top-right frame). On the other hand, the $C2$ set of cuts becomes necessary if one wants to study the properties of all the heavier resonances, as exemplified by comparing in this figure the bottom-left (where only the W_1 resonance is barely visible) to the bottom-right (where the W_2 resonance clearly stands out) plot. In Tab. IV we show the masses and widths of the heavy gauge bosons for the aforementioned 4S benchmark point, so that one can easily identify the origin of the shapes seen in Fig. 5.

Z	M, Γ (GeV)	W	M, Γ (GeV)
1	1012, 33.8	1	1008, 33.3
2	1256, 27.2	2	1255, 26.7

TABLE IV: Numerical values of for masses and widths of the heavy gauge boson resonances for the benchmark point $z = 0.8$, $M_1 = 1$ TeV and $a_{W_1} = 0.13$ of the 4S model.

Other kinematic variables which are efficient to establish the heavy gauge boson resonances of the 4S model are P_M^T and M_{456} , so long that the $C2$ cuts are used. This is exemplified in Fig. 6. Here, one can appreciate the Z_1 and Z_2 resonances stemming in the first variable and the W_1 and W_2 ones emerging from the second observables. One subtlety to observe here though is that in the former case the terms responsible for the effect are the squares of the corresponding 4S diagrams (the peaks in P_M^T are exactly at $M_{Z_1}/2$ and $M_{Z_2}/2$, left plot) whereas in the latter case the relevant contributions are due to the interference of the corresponding 4S graphs with the SM ones (one can appreciate that are the dips and not the peaks which correspond to the W_1 and W_2 masses, right plot). Again, we illustrate such a phenomenology for the case of our usual benchmark scenario, though we can confirm that it remains the same over the entire parameter space of the 4S model.

Furthermore, the lineshapes of the distributions used to extract the gauge resonances (with the $C2$ cuts) of the 4S model are also largely insensitive to the choice between the (a) and (b) solutions in eqs. (5) and (6), respectively. In fact, another very interesting byproduct of the $C2$ cuts is related to the possibility of strictly constraining the model couplings. In fact, as shown in a previous subsection, once all the masses of the model are extracted (thereby z and M_1 are known), two free parameters still remain, a_{W_1} and a_{W_2} . As we can see from Fig. 7 (top-left plot), if the Cc set of cuts is adopted, the interval over which, e.g.,

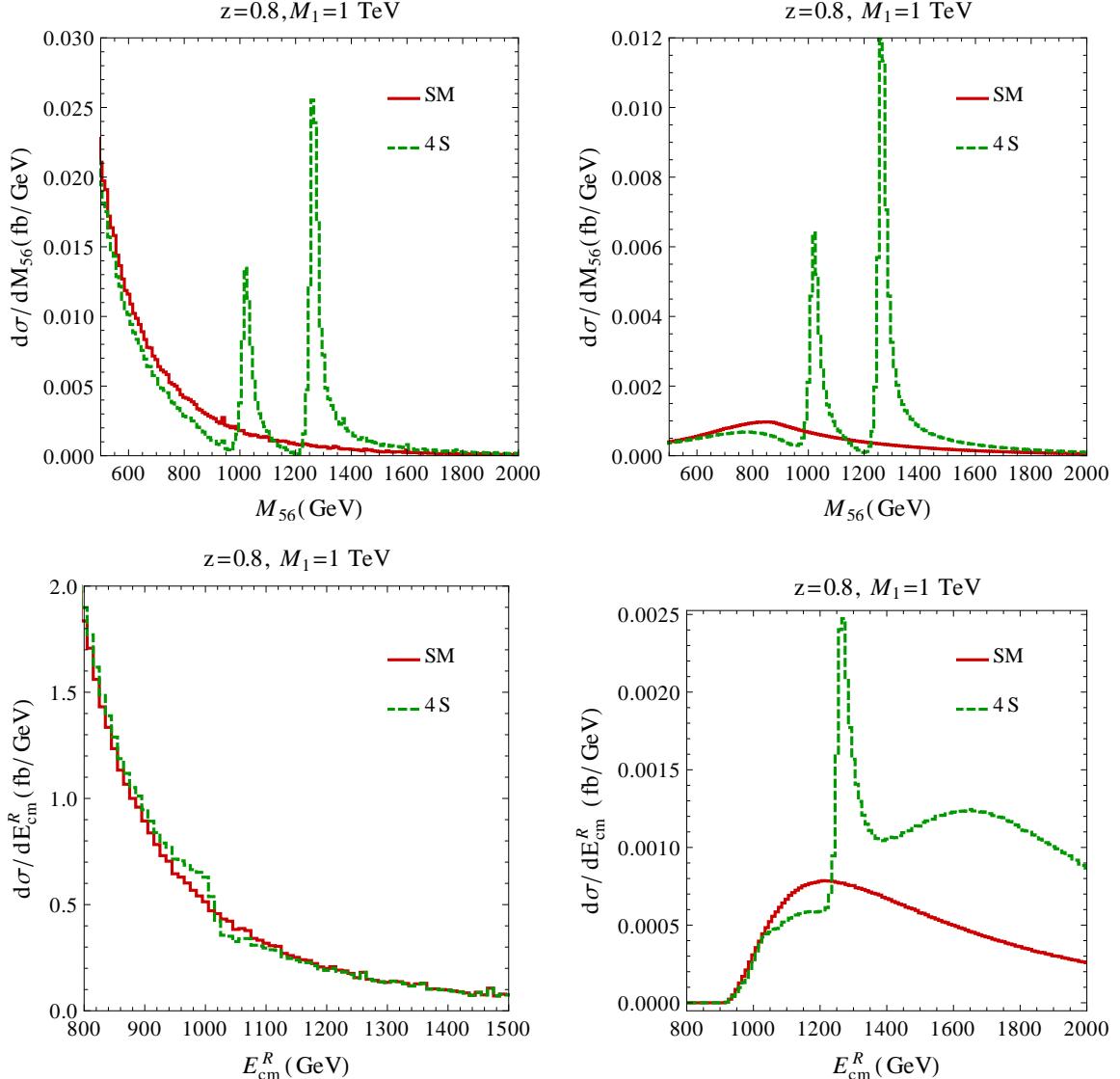


FIG. 5: Differential distributions of (top) the invariant mass of the di-lepton pair coming from the neutral gauge boson (M_{56}) and of (bottom) the reconstructed CM energy at partonic level (E_{cm}^R) for the full 4S model and for the SM, for a benchmark point in parameter space. In particular, we choose $z = 0.8$, $M_1 = 1 \text{ TeV}$ and $a_{W_1} = 0.13$. Left(Right): using the $Cc(C2)$ cuts.

a_{W_1} can be constrained is very large. In fact, it corresponds to the spacings between the points on the blue and green curves intercepting an horizontal line given by a measured CS. Instead, if we consider the $C2$ cuts, for a fixed CS, we have only two very small allowed regions for a_{W_1} , see Fig. 7 (top-right plot) [22]. The same can be said for the WW and DY channels, see Fig. 7 (bottom-left and bottom-right plots, respectively). This is due to the fact that the $C2$ cuts are very sensitive to a_{W_1} (but not to a_{W_2}) so that, once we have

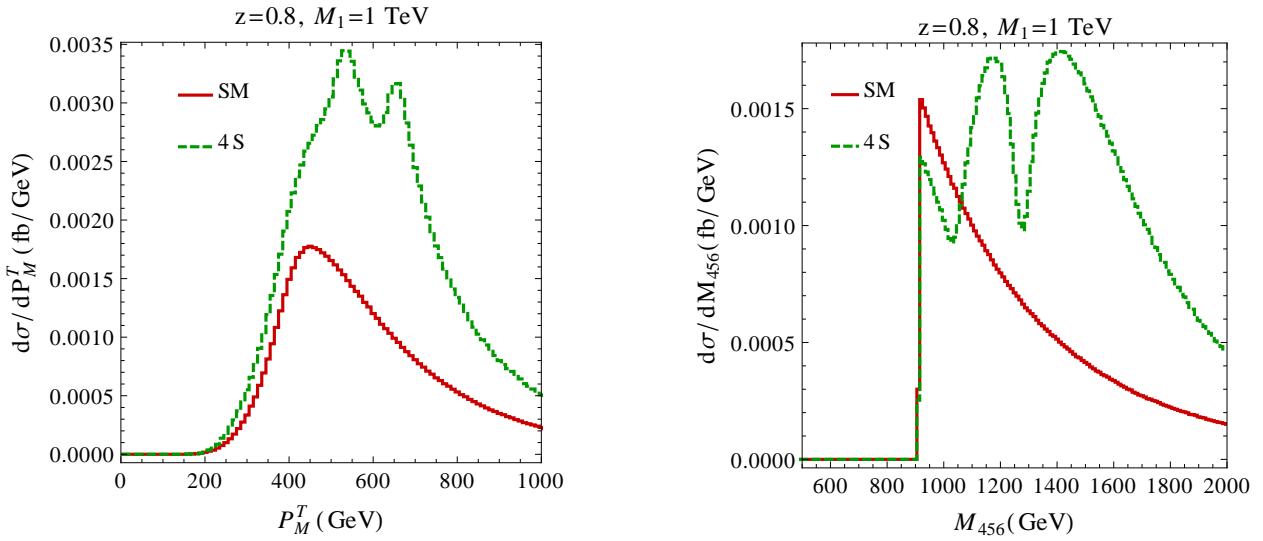


FIG. 6: Differential distributions of (left) the transverse momentum of the di-lepton pair coming from the neutral gauge boson (P_M^T) and of (right) the invariant mass of the three leptons (M_{456}) for the full 4S model and for the SM one, for a benchmark point in parameter space. In particular, we choose $z = 0.8$, $M_1=1$ TeV and $a_{W_1} = 0.13$. Here, we are using $C2$ cuts.

a precise measurement of a_{W_1} , we can re-use the sample defined by Cc cuts (that are very sensitive to a_{W_2} but not to a_{W_1}) in order to extract the appropriate value of a_{W_2} . Further notice that, despite, for consistency with the WW case studied in Ref. [4], we have limited the exercise here for the WZ process to the case of the $e\mu$ flavour combination only, this can equally be done for the $3l$ one. Finally, the same procedure can be applied to any point on the 4S parameter space, although we have illustrated it here for our usual benchmark point.

In closing then, we should conclude that a judicious use of the two sets of cuts introduced can enable one to extract efficiently both the mass and coupling spectrum of the 4S model. However, the spectrum reconstruction described here requires rather large luminosities, so that, whilst it is certainly applicable to the LHC at 14 TeV, it remains of limited scope at 7 and 8 TeV.

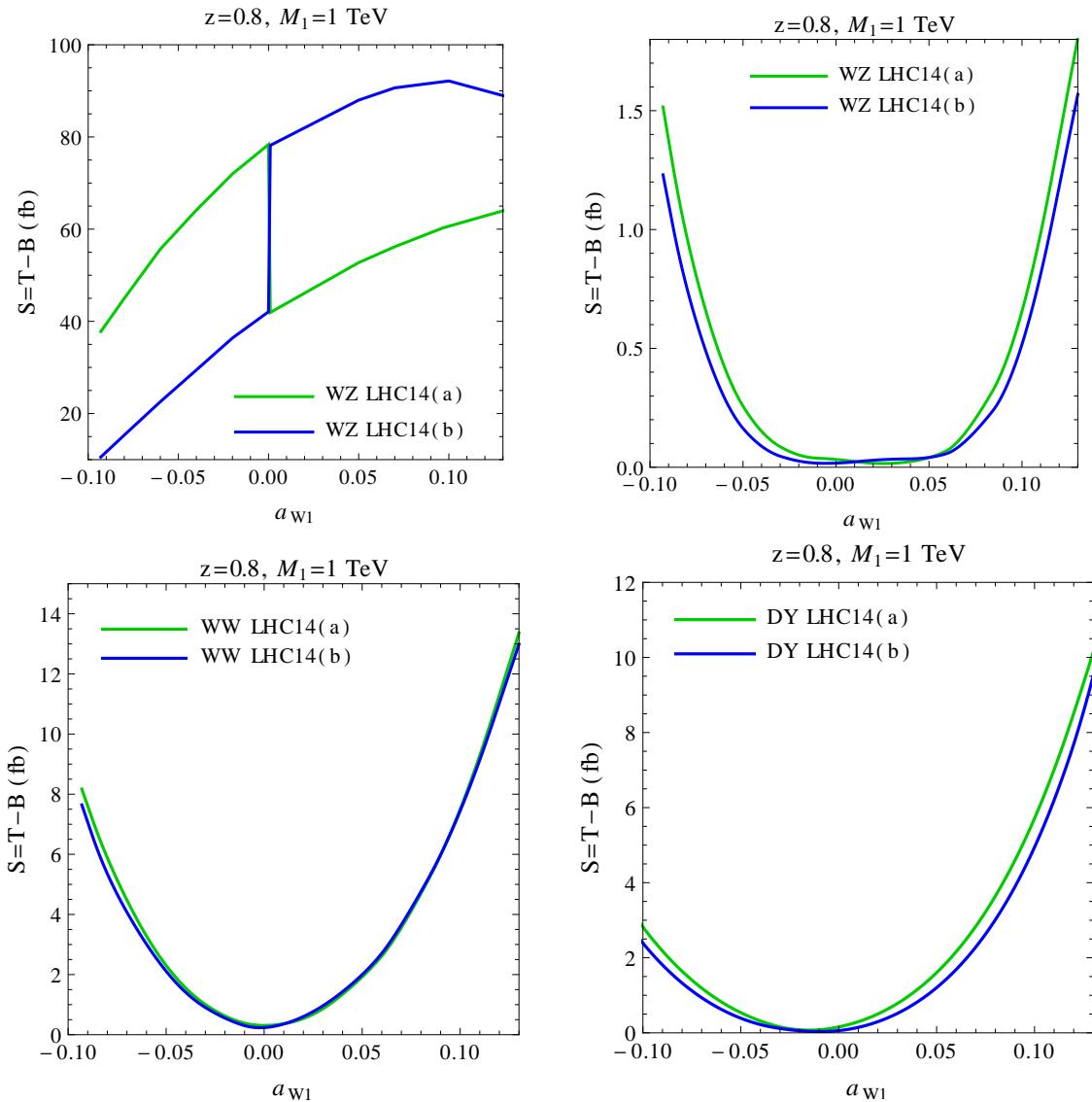


FIG. 7: Correlations between the CS for the WZ signal process, defined as the difference between the full 4S result and the SM one, and the coupling a_{W1} for the benchmark point $z = 0.8$ and $M_1 = 1$ TeV. Top-left(right): using $Cc(C2)$ cuts. Bottom-left(right): the same for WW (DY) using the So (standard) cuts introduced in [4]([7, 9]). The labels (a) and (b) refer to the solutions in eqs. (5) and (6), respectively. These values refers to the $e\mu$ process.

IV. CONCLUSIONS AND OUTLOOKS

In summary, mixed di-boson production (via WZ diagrams) in a 4S model supplemented by a composite Higgs state revealed itself as an important LHC channel in order to test such a scenario of EWSB. On the one hand, it crucially contributes to the discovery potential of the LHC over new regions of the parameter space, if compared to the scope of the twin charged di-boson (via WW topologies) mode and DY processes, specifically, for high gauge boson

masses. On the other hand, thanks to the fact that the source of missing (transverse) energy in the case of the WZ channel is due to only one neutrino (unlike the case of WW mediation, where it is due to two neutrinos), one can reconstruct fully the final state kinematics, hence extracting all the intervening resonances, in turn implying a complete knowledge of the gauge boson spectrum of the 4S model. This can be achieved after the sequential implementation, firstly, of looser cuts merely emulating detector acceptance (that however extract the typical mass scale of the lowest lying new resonances) and, secondly, of a tighter selection exploiting such kinematical information (that enables an effective mass spectroscopy). Such conclusions are generally valid independently of the LHC setup, though they become quantitatively most relevant at higher energies and luminosities.

Finally, we should emphasise that a byproduct of our analysis was the realisation that the aforementioned looser cuts are primarily sensitive to interference effects (between the genuine 4S contributions and the SM ones) whereas the tighter ones enhance instead resonance effects (primarily of the new heavy gauge bosons). This therefore calls, in the same spirit as in Ref. [10] in the DY case, for revisiting the scope of the di-boson channels (both charged and mixed) at the LHC, under assumptions different from those routinely made by the LHC experiments (of new resonance dominance). Also, one other aspect so far unexplored of processes (1) and (3) which does not pertain to the channels in (2) is the dependence of the CS on tri-linear gauge couplings, which will also constitute the subject of another publication.

Acknowledgments

We all thank Elena Accomando for discussions. SM is financed in part through the NExT Institute. LF thanks Fondazione Della Riccia for financial support. The work of SDC and DD is partly supported by the Italian Ministero dell’Istruzione, dell’Università e della Ricerca (MIUR) under the COFIN programme (PRIN 2008).

- [1] E. Accomando, S. De Curtis, D. Dominici, and L. Fedeli, Phys. Rev. **D79**, 055020 (2009), 0807.5051.
- [2] G. Aad et al. (ATLAS Collaboration), Phys. Lett. **B716**, 1 (2012), 1207.7214.

- [3] S. Chatrchyan et al. (CMS Collaboration) (2012), 1207.7235.
- [4] E. Accomando, L. Fedeli, S. Moretti, S. De Curtis, and D. Dominici (2012), 1208.0268.
- [5] E. Accomando, S. De Curtis, D. Dominici, and L. Fedeli, Nuovo Cim. **123B**, 809 (2008), 0807.2951.
- [6] E. Accomando, A. Belyaev, L. Fedeli, S. F. King, and C. Shepherd-Themistocleous, Phys. Rev. **D83**, 075012 (2011), 1010.6058.
- [7] E. Accomando, S. De Curtis, D. Dominici, and L. Fedeli, Phys. Rev. **D83**, 015012 (2011), 1010.0171.
- [8] E. Accomando, D. Becciolini, S. De Curtis, D. Dominici, and L. Fedeli (2011), 1105.3896.
- [9] E. Accomando, D. Becciolini, S. De Curtis, D. Dominici, and L. Fedeli (2011), 1107.4087.
- [10] E. Accomando, D. Becciolini, S. De Curtis, D. Dominici, L. Fedeli, et al., Phys. Rev. **D85**, 115017 (2012), 1110.0713.
- [11] H. Murayama, I. Watanabe, and K. Hagiwara (1992).
- [12] T. Stelzer and W. Long, Comput. Phys. Commun. **81**, 357 (1994), hep-ph/9401258.
- [13] G. P. Lepage, J. Comput. Phys. **27**, 192 (1978).
- [14] H. Lai et al. (CTEQ Collaboration), Eur. Phys. J. **C12**, 375 (2000), hep-ph/9903282.
- [15] F. Bach and T. Ohl, Phys. Rev. **D85**, 015002 (2012), 1111.1551.
- [16] Notice that $z = 0.8$ is the case for which processes (2) and (3) have comparable scope at the LHC, as seen in Ref. [4].
- [17] Notice that, in performing this exercise, in order to obtain a finite result, we ought to keep a finite mass for the muon, which we take to be $m_\mu = 0.105$ GeV.
- [18] As choice (a) was made in Refs. [4, 7, 9], to which we will compare, we will maintain it here as our default, though we will comment regarding the alternative choice.
- [19] Notice that di-lepton final states with identical flavours are of no use for the WW process, as they are burdened by an overwhelming SM background induced by ZZ events, with one Z boson decaying invisibly.
- [20] Hereafter, we choose to use a 10 GeV bin for the differential distributions due to the fact that the experimental resolution (in mass and transverse momentum) is of order 1% at 1–2 TeV for the electron, while for the muon the rate is more like 5%.
- [21] In contrast, note that the more kinematically constrained structure of DY processes does not allow one to apply the same algorithm in the case of the charged current channel in (2), so

that here the extraction of the masses and couplings of the 4S charged gauge bosons is only possible through a transverse mass [9], notoriously less sensitive to the actual resonant mass than an invariant mass distribution.

- [22] This however occurs at the cost of a much reduced cross section, hence of a larger statistical error on the corresponding experimental measurement, so that a compromise between the Cc and $C2$ cuts may have to be found.